Quantum Transport Devices Based on Resonant Tunneling

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Introduction

 Some general aspects of resonant tunneling diodes will be discussed
 RTDs can be considered as devices which are in active competition with conventional CMOS

> Transfer Matrix Method

- Electrons have a wave like character
- In structures with dimensions in the range of electron wavelength, quantum mechanical transport becomes relevant
- One of these transport mechanisms is the tunneling process
- Electrons can penetrate through and traverse a potential barrier with a finite transmission probability independent of temperature
- In classical view electrons can overcome a potential barrier only thermodynamically

- Envelope function description of electron state : rapid changing electron potential is approximated with an envelope potential
- Envelope function is based on effective mass description of the band structure and leads to electron effective mass Schrödinger equation :

$$\left[-\frac{\overline{h}^2}{2}\frac{d}{dz}\frac{1}{m^*(z)}\frac{d}{d(z)}+\Phi(z)\right]\Psi(z)=W_z\Psi(z)$$

 $\Psi(z)$: electron wave function, W_z : electron energy is Z direction m^{*}: electron effective mass, $\Phi(z)$: potential energy at the conduction band minimum

- Occupation probabilities can be predicted from absolute square of wave function $|\Psi(z)|^2$
- Consider a sequence of n different layers (fig 1) with different potential energies (φ_i) and electron effective masses (m^{*}_i)



Figure 1: Sequence of n different layers; in each layer the effective mass m_i^* and the potential Φ_i was assumed to be constant.

 Ψ_i(z) can be written as a superposition of propagating waves in z and –z direction with amplitudes A_i and B_i

$$\Psi_i(z) = A_i e^{ik_i z} + B_i e^{-ik_i z} \coloneqq A_i \Psi_{i+} + B_i \Psi_{i-}$$

boundary conditions :

$$\Psi_{i}(z_{i}) = \Psi_{i+1}(z_{i})$$
$$\frac{1}{m_{i}^{*}} \frac{d}{dz} \Psi_{i}(z_{i}) = \frac{1}{m_{i+1}^{*}} \frac{d}{dz} \Psi_{i+1}(z_{i})$$

In matrix form :

$$definition: TM_{i} \coloneqq \begin{bmatrix} \Psi_{i+} & \Psi_{i-} \\ \frac{1}{m_{i}} \Psi_{i+}^{'} & \frac{1}{m_{i}} \Psi_{i-}^{'} \end{bmatrix}$$
$$TM_{1(z=z1)} \begin{bmatrix} A_{1} \\ B_{1} \end{bmatrix} = TM_{2(z=z1)} \begin{bmatrix} A_{2} \\ B_{2} \end{bmatrix}, TM_{2(z=z2)} \begin{bmatrix} A_{2} \\ B_{2} \end{bmatrix} = TM_{3(z=z2)} \begin{bmatrix} A_{3} \\ B_{3} \end{bmatrix}, \dots$$
$$TM_{n-1(z=z_{n-1})} \begin{bmatrix} A_{n-1} \\ B_{n-1} \end{bmatrix} = TM_{n(z=z_{n-1})} \begin{bmatrix} A_{n} \\ B_{n} \end{bmatrix},$$

 Amplitudes of the propagating waves in z and –z direction in last layer can be written as:

$$\begin{bmatrix} A_n \\ B_n \end{bmatrix} = TM \begin{bmatrix} A_1 \\ B_1 \end{bmatrix}$$

 $TM = TM_{n(z=z_{n-1})}^{-1} ... TM_{2(z=z^2)} TM_{2(z=z^2)}^{-1} TM_{1(z=z^1)}$

• Transmission probability Tc can be written as the ratio of outgoing to the incoming quantum mechanical probability current: $k m_{*}^{*} |A|^{2} \det TM$

$$T_{c} = \frac{k_{n}}{k_{1}} \frac{m_{1}^{+}}{m_{n}^{*}} \frac{|A_{n}|^{2}}{|A_{1}|^{2}}, A_{n} = \frac{\det TM}{TM_{22}} A$$
$$\det TM = \frac{k_{1}m_{n}}{k_{n}m_{1}}$$
$$T_{c} = \frac{k_{1}}{k_{n}} \frac{m_{n}^{*}}{m_{1}^{*}} \frac{1}{|TM_{22}|^{2}}$$

Tunneling through a single barrier A single potential barrier is shown in figure (fig2) AlAs barrier embedded in GaAs

Figure 2: Schematic band diagram of a single AlAs barrier (a) and the corresponding tunneling transmission probability for different barrier thicknesses (b).



- Transmission probability is calculated as a function of electron energy
- Finite transmission probability for electrons below potential height of 1 eV (tunneling)
- The smaller the barrier thickness the higher the tunneling probability

- Tunneling through a double barrier structure
 - Figure (fig 3) shows the case of tunneling through a double barrier structure
 - 4 nm tick AIAs barriers separated by a 5 nm GaAs well

Figure 3: Schematic band diagram of a double barrier structure of AlAs embedded in GaAs (a) and the corresponding tunneling transmission probability (b).



- In contrast to single barrier, there are three sharp maxima below 1 eV
 - Interpreted as quasi-bound states with narrow energetic bandwidth through which electrons can tunnel through open channels in the barrier
- This is not describable by a sequential picture of two wells
- Quantum mechanical devices cannot be put too close together without changing the characteristics of the single device

Resonance properties

- Resonant tunneling diode is the experimental realization of double barrier structure
- Figure (fig 5) shows the behavior of resonances
- A resonance can be considered as a channel which opens electron flux, current density first increases then decreases



Current voltage characteristics

- Current density can be calculated based on transmission probability and the corresponding occupation densities
- Text gives a relation for calculating current density based on the potential profile φ of the structure
- The potential can be obtained by coupling effective mass Schrödinger equation with Poisson equation in a self-consistent manner

- Figure (fig 6) shows a typical currentvoltage characteristic
- Negative differential resistance is a main feature



 The quantum device simulation package NEMO (NanoElectronic MOdeling) simulates a wide variety of quantum devices including RTDs

Interface and growth temperature

(fig 8)

- PVR (Peak to Valley Ratio) is a merit of quality for RTDs
- Highest PVR coincides with sharpest interface between barriers and the well
- Temperature range between 580 and 600 results in highest PVR for AIAs/GaAs RTDs



> Operation Speed of RTDs

- One of the most attractive features of RTDs is their potential for extremely high speed operation
- RTDs with 712 GHz oscillation and 1.5 ps switching times have been reported
- It is important to differentiate "tunneling time" and "RC time"
- Tunneling time is in order of the resonant-state lifetime or escape time which is the time it takes an electron in the quantum well to escape from it:

- Shorter tunneling times can be obtained with thinner and lower barries
- Various non-idealities affect tunneling time in real RTDs
- Tunneling time determines the intrinsic delay of RTDs

- In most applications, operation speed of RTDs is limited not by the intrinsic tunneling time but by the charging time of RTD capacitance
- Equivalent circuit of RTDs is shown in figure (fig 10)
- The capacitance-voltage curve is also shown in the figure
 (a)



> Applications of RTDs

- Several applications exploit negative differantial resistance (NDR) of RTDs
- Resonant tunneling transistors
 - To make a three terminal tunneling device RTDs are merged with conventional transistors and resonant tunneling bipolar transistors, resonant tunneling hot electron transistors and gated RTDs are fabricated

- Gated RTDs have Schottkey or junction gates around the emitter to control RTD area
- Concept of Monostable-Bistable transition logic elements (MOBILES)
 - The ultrahigh-speed logic gate : MOBILE exploits NDR of the RTDs
 - A circuit consisting of two NDR devices connected serially, to exploit monostable to bistable transition
 - Bias voltage is oscillated to generate the transition
 - NDR devices with third terminal are used to modulate their peak currents

- Figure (fig 12) shows the load curves and the corresponding potential energy diagrams
- Figure (fig 13) shows the operation of a simple



Figure 13: Circuit configuration and the operation wave form of the simple MOBILE inverter. Traces are the clock, input, and output voltages, from top to bottom, respectively.



Integration Technology for MOBILES

- NDR devices with third terminal are required for MOBILEs
- Gated RTDs were first used as three terminal NDR device, however they have disadvantages including high capacitance, difficult to optimize the layer structure, low PVR and difficult to fabricate
- To overcome the problems, RTDs connected in parallel with HEMTs were fabricated

Resonant Tunneling Devices > Examples of MOBILEs

• Figure (fig 17, 18) shows implementation of a weighted sum threshold logic function



Figure 17: The circuit configuration and the operation of three-input MOBILE. The traces are the input 1 (w = 1), the input 2 (w = 2), the input 3 (w = 4), and the output, from top to bottom, respectively. The weighted sum (S) is shown at the bottom of the figure as a reference. The threshold value is selected by the control voltage, V_{con} , and it is chosen to be 2.5 in this figure.



Resonant Tunneling Devices
 Figure (19, 20, 21) demonstrate a test circuit for evaluating operation speed of the MOBILE and the estimated power consumption

Figure 20: The operation waveform of the MOBILE D-FF (a)

(b)

(a) with the input of a pseudo-random bit stream at 35 Gb/s,

(b) with an input bit pattern of (0111010000101110).(The upper trace in b) is the complement of the input data stream.)



- One of the promising applications of MOBILEs is analog-to-digital-converter (ADC)
- Figure (fig 22) shows a block diagram of a ΔΣ ADC



- ΔΣ modulator converts analog input into a pulse density at a frequency much higher than the Nyquist rate
- The filter cuts the high frequency component and down converts pulse density into the high-resolution digital output at Nyquist rate
- Higher resolution can be obtained by increasing sampling rate
- This method does not require a high accuracy analog component

- MOBILEs can be used to fabricate high performance $\Delta\Sigma$ modulator
- Adder and shifter circuits used in digital filter can also be fabricated by MOBILEs
- Figure (fig 23) shows a ΔΣ modulator based on MOBILE

